



Proceedings of the Eighteenth International Conference on  
Civil, Structural and Environmental Engineering Computing  
Edited by: P. Iványi, J. Kruis and B.H.V. Topping  
Civil-Comp Conferences, Volume 10, Paper 15.1  
Civil-Comp Press, Edinburgh, United Kingdom, 2025  
ISSN: 2753-3239, doi: 10.4203/ccc.10.15.1  
©Civil-Comp Ltd, Edinburgh, UK, 2025

# **A Less Complex Wing Theory in Ideal Fluids**

**R. W. Meyer**

**Department of Mechanical Engineering and Maritime Studies,  
Western Norway University of Applied Sciences,  
Haugesund, Norway**

## **Abstract**

In this paper the author suggests a less complex wing theory in ideal fluids that connects lift force and drag force of a finite 3-dimensional wing with mass flow into the generated vortex cores.

Suggesting an induced drag in ideal fluids for 2-dimensional wing sections, leads to a modified version of the equation Prandtl has derived for the relation between lift force  $F_y$  and drag force  $F_x$  of a 3-dimensional finite wings. This new equation goes over into Prandtl's equation, if we have low aspect ratios of the wings.

The author derives equations that calculates the vortex core radius of the generated vortexes and are connecting the mass flow into the vortex cores with drag and lift of the wing. Additionally, the author shows, that the power to compensate for the drag of a plane in an ideal fluid is consumed by the rotational power of the generated vortex cores.

**Keywords:** ideal fluids, wing theory, vortex core, lift coefficient, drag coefficient, random walks.

## **1 Introduction**

A model for unsteady potential flow in ideal fluids, called the Random Walk Source Model is presented in [1]. This model generates the shape of a particle front emerging from a line source, creeping towards and past a wing section. The lift coefficient can be predicted by the difference of creeping distance between the two front segments that have been separated by the wing section and an extended separation line which

maintains the Kutta condition. The front is a result of a model of unsteady potential flow, and further leads to an unexpected developments of an equation leading to the formulation of an induced drag in 2- dimensions in ideal fluids. We have to account for this part of the drag also in 3- dimensions, which is made in the following description.

Wake vortices behind planes are a typical topic in fluid dynamics. Lanchester [2] and Prandtl [3] described the 3-dimesional flow behind wings with finite span in ideal fluids, resulting in an induced drag out of the ability to generate a circulation, in a relation between lift force  $F_y$  and expected drag force  $F_x$ :

$$F_x = F_y \cdot \frac{w_\infty}{2u_\infty} \quad (1)$$

Where  $u_\infty$  is the velocity of the plane and  $w_\infty$  is the induced downward velocity of the fluid flow.

As described in [4] the relation between  $C_l$  and  $C_d$  is stated as an induced drag in 3-dimensional flow in ideal fluids and now additionally an induced drag in 2-dimensional flow:

$$C_d = \left(\frac{C_l}{2\pi}\right)^2 \cdot \frac{\pi \cdot \ln(1+\sqrt{2})}{\sqrt{8}} + \frac{C_l^2}{\pi\Lambda} \quad (2)$$

This is leading to the equation:

$$F_x = F_y \cdot \frac{w_\infty}{2u_\infty} \cdot \left[ \frac{\Lambda \cdot \ln(1+\sqrt{2})}{4 \cdot \sqrt{8}} + 1 \right] \quad (3)$$

Where  $\Lambda$  is the aspect ratio of the wing. This is an extension of formular (1). Formular (2) is gradually going over into formular (1) when  $\Lambda$  is close or less 1.

## 2 Level flight of a wing of finite span in ideal fluids

We calculate  $C_l$  out of the lift force  $F_y$  and the speed  $u_\infty$  of the plane together with the density  $\rho$  and the area  $A$  of the wing:

$$C_l = \frac{2 \cdot F_y}{\rho \cdot u_\infty \cdot A} \quad (4)$$

$C_d$  the drag coefficient of the wing is then calculated according to equation (2).

Alternatively the author suggests to formulate the drag coefficient with help of the mass flow  $\dot{m}_{core}$  into the vortex cores generated from the wing and the descent velocity  $w_\infty$ .

$$C_d = \frac{2 \cdot \dot{m}_{core} \cdot w_\infty}{\rho \cdot u_\infty^2 \cdot A} \quad (5)$$

The mass flow  $\dot{m}_{core}$  is defined as

$$\dot{m}_{core} = C_d \cdot \frac{\rho}{2} \cdot A \cdot \frac{u_\infty^2}{w_\infty} \quad (6)$$

and  $w_{oo}$  is defined as

$$w_\infty = \frac{\Gamma_0}{b_1} \quad (7)$$

Where  $b_1$  is the wingspan, and  $\Gamma_0$  the circulation

$$\Gamma_0 = 4 \cdot T_1 \cdot u_{HK} \quad (8)$$

Where  $T_1$  is the average wing width, the wing area is  $A_1 = b_1 \cdot T_1$  with

$$u_{HK} = \frac{C_l u_\infty}{2 \cdot \pi} \quad (9)$$

We can then write the drag force

$$F_x = \dot{m}_{core} \cdot w_\infty \quad (10)$$

And the lift force

$$F_y = \dot{m}_{core} \cdot \frac{C_l}{C_d} \cdot w_\infty \quad (11)$$

The vertical mass flow generating the lift force is according to Prandtl the flow through a circle area with the wing span as the diameter. Fluid particles are entering this area with  $u_\infty$ , perpendicular to this area. This mass flow is moving down with  $w_\infty$ , parallel to the area

$$F_y = \dot{m}_y \cdot w_\infty \quad (12)$$

Where  $\dot{m}_y$  is defined (as explained above)

$$\dot{m}_y = \dot{V}_1 \cdot \rho \quad (13)$$

with

$$\dot{V}_1 = A_{01} \cdot u_\infty \quad (14)$$

Here  $A_{01}$  is a circle with the wing span as the diameter

$$A_{01} = \frac{\pi \cdot b_1^2}{4} \quad (15)$$

The relation between  $\dot{m}_y$  and  $\dot{m}_{core}$  is

$$\frac{\dot{m}_y}{\dot{m}_{core}} = \frac{C_l}{C_d} \quad (16)$$

### 3 Comparing calculated core radius with published experiments

In Figure 1 the calculated vortex core

$$r_{core} = \sqrt{\frac{A}{4} \cdot \frac{C_l}{2\pi}} \quad (17)$$

is compared to an experiment published in [5]. The calculated tangential velocity and the core radius is fitting good into the experimental data.

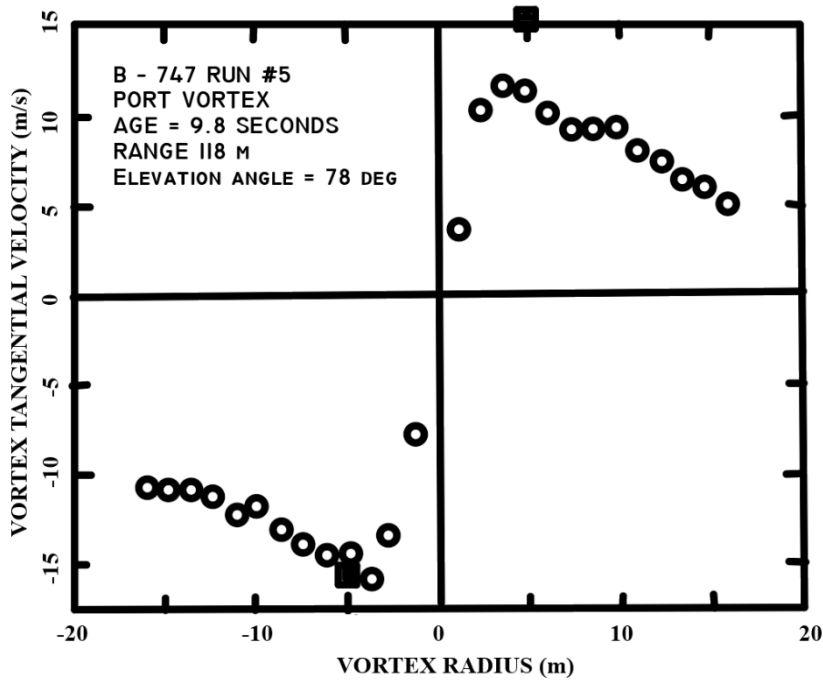


Figure 1: Vortex radius and vortex tangential velocity reproduction from [5] circles are experimental data, squares ■ are calculated according (17)

$$r_{core} = \sqrt{\frac{A}{4} \cdot \frac{C_l}{2\pi}} = 4.99 \text{ m}$$

and  $u_{HK} = \frac{C_l u_\infty}{2 \cdot \pi} = 15.26 \text{ m/s}$  the vortex tangential velocity.  $C_l = 1.23$  is the lift coefficient,  $u_\infty$  is the velocity of the plane here 78 m/sec and A is the wing area 510 m<sup>2</sup>

## 4 Power calculation

Here we write down the power that is consumed to keep the wing flying on level consisting of the drag force multiplied with the velocity of the wing

$$P^*_{tot} = F_x \cdot u_\infty \quad (18)$$

Alternatively we can see on the descent velocity  $v_y$  of the unpowered wing in relation to the horizontal velocity  $u_{oo}$  of the wing with the known relation between  $C_d$  and  $C_l$ . We use the relation

$$\frac{C_d}{C_l} = \frac{v_y}{u_\infty} \quad (19)$$

That leads to the descent velocity

$$v_y = \frac{C_d \cdot u_\infty}{C_l} \quad (20)$$

The lifting force multiplied with the descent velocity is the power to keep the plane on level according to lift and drag.

$$P^*_{tot} = F_y \cdot v_y = \dot{m}_y \cdot w_\infty \cdot v_y \quad (21)$$

The kinetical energy of the vortex cores (seen as cylinders with radius  $r_{core}$ ) is calculated

$$E_{kin} = \frac{1}{2} \cdot \Theta \cdot \omega^2 \quad (22)$$

Where  $\Theta$  is the momentum of inertia of a cylinder (the vortex core) which is growing in length in time leading to  $\dot{\Theta}$

The rotational power of the vortex cores (kinetical energy per time) is calculated

$$P^*_{core} = \dot{E}_{kin} = \frac{1}{2} \cdot \dot{\Theta} \cdot \omega^2 \quad (23)$$

Where  $\omega$  is the angular velocity

$$\omega = \frac{u_{HK}}{r_{core}} \quad (24)$$

Inertial mass  $\Theta$  of a cylinder (representing the vortex cores)

$$\Theta = \frac{1}{2} m_{core} \cdot r_{core}^2 \quad (25)$$

Ongoing flow into the vortex core (the vortex core is a cylinder with constant radius  $r_{core}$  increasing the mass with growing length)

$$\dot{\Theta} = \frac{1}{2} \dot{m}_{core} \cdot r_{core}^2 \quad (26)$$

Leading to the rotational power of the vortex cores

$$P_{core}^* = \frac{1}{4} \cdot \dot{m}_{core} \cdot u_{HK}^2 \quad (27)$$

Leading to the relation

$$\frac{P_{tot}^*}{P_{core}^*} = \frac{C_l}{C_d} \quad (28)$$

### 3 Results

Figure 2 is showing a survey over the suggested equations.

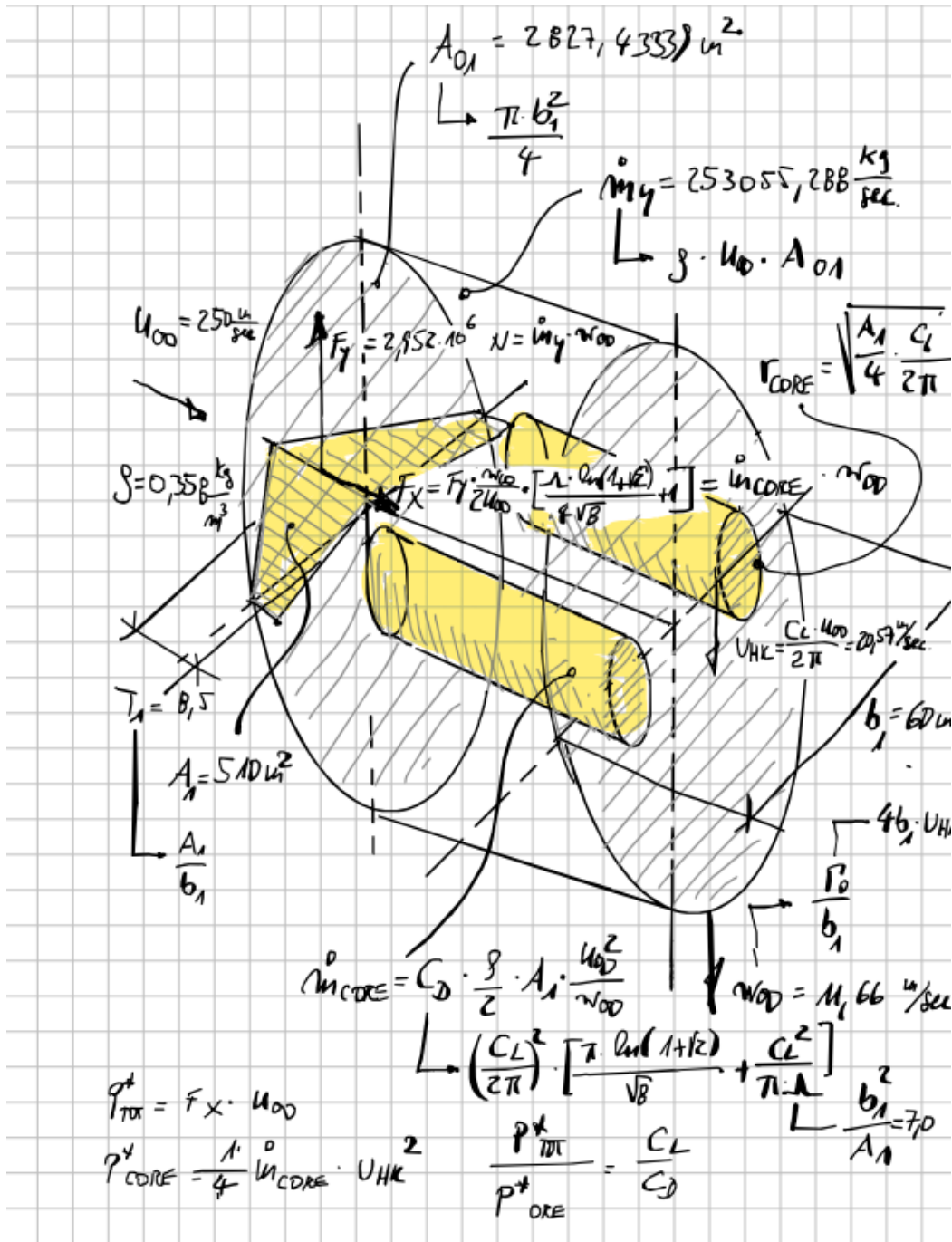


Figure 2: Visualization of lift and drag of a Boeing 747 according to the theory suggested in this paper. Plane parameter from [6] calculated values in Table 1

Description	Parameter	Value	
Lift force	$F_y$	2952000 N	
Air density	$\rho$	0.358 kg/m <sup>3</sup>	
Plane velocity	$u_\infty$	250 m/s	
Wing span	$b_1$	60 m	
Wing area	$A_1$	510 m <sup>2</sup>	
Circulation	$\Gamma$	$699.92 \frac{m^2}{s}$	
Aspect ratio	$\Lambda$	7.058	
	Calculated values		Used equation
Drag force	$F_x$	106746,16 N	(3)
Mass descent velocity	$w_\infty$	11,6654 m/s	(7)
Area of influenced mass	$A_{01}$	2827,433 m <sup>2</sup>	(15)
Downward accelerated volume	$\dot{V}_1$	$706858.347 \frac{m^3}{s}$	(14)
Lift coefficient	$C_l$	0.517384	(4)
Average wing section	$T_1$	8.5 m	
Trailing edge velocity	$u_{HK}$	20.586 m/s	(9)
Drag coefficient	$C_d$	0.01870893	(2)
Mass flux into the cores	$\dot{m}_{core}$	$9150.6369 \frac{m^3}{s}$	(6)
Mass flow vertical	$\dot{m}_y$	$253055.2882 \frac{m^3}{s}$	(13)
Core radius	$r_{core}$	3.240199 m	(17)
Total power used	$P^*_{tot}$	26686539.8 Nm/s	(18)
Power consumed from the vortex cores	$P^*_{core}$	969477.8 Nm/s	(27)
Power relation	$\frac{P^*_{tot}}{P^*_{core}} = \frac{C_l}{C_d}$	27.65	(28)

Table 1: Input parameter Boeing 747 [6] and from the author calculated values.



## 4 Conclusions

A wing section has no drag in ideal fluids according to the potential theory. We can calculate the circulation which is leading to the lift force. The start vortex is moved to vicinity and has lost the influence in the theoretical model. In reality we have a vortex pair in 2- dimensions which is continuously fed with attracted fluid particles. In the random walk source model [1],[4] it is shown, that you can identify an induced drag from wing sections in ideal fluids in 2- dimensions which is resulting in an equation that connects lift coefficient and drag coefficient. This is unexpected. In [4] it is shown that you can calculate the relation of this coefficients for a rotating cylinder with endplates. Also this is unexpected because we assume, that we have a highly turbulent flow behind the cylinder. In this paper the author transfers the knowledge of induced drag in 2- dimensions to wings in ideal fluids in 3- dimensions, se Figure 3. It is shown, that the core radius of vortex cores is depending on the lift coefficient and the wing area. Attracted fluid particles are feeding the vortex core generating the drag through the rotational power in the vortex cores. It is presented an equation (3) which is a modification of Prandl's relation (1) between lift force and drag force of wings in ideal fluids. In the presented equations the relation between lift coefficient and drag coefficient is the same relation as between the mass that is moved downward from the plane (to generate the lift) compared to the mass that is flowing into the vortex cores (16). The total power to keep the plane on level (18) divided by the power to keep the vortex cores rotating (27) is  $\frac{C_l}{C_d}$  (28).

$$\begin{aligned}
 & \dot{V}_1 = A_{01} \cdot u_\infty \\
 & \dot{m}_y = \dot{V}_1 \cdot \rho \\
 & \dot{m}_{core} = \dot{m}_y \cdot \frac{C_l}{C_d} \\
 & \dot{m}_{core} = C_d \cdot \frac{\rho}{2} \cdot A \cdot \frac{u_\infty^2}{w_\infty} \\
 & C_l = 2 \cdot \pi \cdot \frac{u_{HK}}{u_\infty} \\
 & F_y = C_l \cdot \frac{\rho}{2} \cdot u_\infty^2 \cdot A \\
 & F_x = F_y \cdot \frac{w_\infty}{2u_\infty} \cdot \left[ \frac{\Lambda \cdot \ln(1+\sqrt{2})}{4 \cdot \sqrt{8}} + 1 \right] \approx F_x = \dot{m}_{core} \cdot w_\infty \\
 & F_x = C_d \cdot \frac{\rho}{2} \cdot u_\infty^2 \cdot A \\
 & r_{core} = \sqrt{\frac{A}{4} \cdot \frac{u_{HK}}{u_\infty}} \\
 & w_\infty = \frac{2 \cdot C_l \cdot u_\infty}{\pi \cdot \Lambda} \\
 & C_d = \left( \frac{C_l}{2\pi} \right)^2 \cdot \frac{\pi \cdot \ln(1+\sqrt{2})}{\sqrt{8}} + \frac{C_l^2}{\pi \Lambda} \\
 & \dot{m}_{core} = 2 \cdot \pi \cdot r_{core}^2 \cdot \rho \cdot u_\infty \cdot \left[ \frac{\Lambda \cdot \ln(1+\sqrt{2})}{4 \cdot \sqrt{8}} + 1 \right]
 \end{aligned}$$

Figure 3: Equations connected with arrows, showing the calculation from generated lift force to drag force of a finite wing in ideal fluids. Calculation of lift and drag of a Boeing 747 according to the theory suggested in this paper taking into account the volume flux into the vortex cores and the vertical vortex flux generated by the plane. Plane parameter from [6] calculated values explained in Table 1.

## Acknowledgements

This work is funded by the Western Norway University of Applied Sciences, Haugesund, Norway, by funding the research group “Random walk source model in fluid dynamics”.

## References

- [1] R. Meyer, “Random walks and hydrodynamical lift from wing sections”, *Physica A*, 242(1–2), 230–238, 1997. [https://doi.org/10.1016/S0378-4371\(97\)00206-9](https://doi.org/10.1016/S0378-4371(97)00206-9)
- [2] F. W. Lanchester, “Aerodynamik Bd. I und Bd. II“, Berlin und Leipzig, 1909 und 1911.
- [3] L. Prandtl, “Tragflügeltheorie, I. und II. “Mitteilungen und Nachrichten der Kgl. Ges. Wiss. Göttingen, Math.- Phys. Klasse S. 451 -477 , 1918 und S. 107 -137”, 1919.
- [4] R. W. Meyer, S. Erland, “Induced drag in two dimensions in ideal fluids”, *Journal of Physics Communications*, Volume 3, Number 11 2019 <https://doi.org/https://doi.org/10.1088/2399-6528/ab5022>
- [5] D. C. Burnham, J. N. Hallock, I. H. Brashears, M. R. Barber, “Ground- Based Measurements of the Wake Vortex Characteristics of a B747 Aircraft in Various Configurations”, FAA- RD-78-146, 1978.
- [6] T. Ehret, “Numerische Simulation und stabilitätstheoretische Untersuchung von Flugzeug- Nachlaufwirbeln“, Dissertation, Fakultät für Maschinenbau der Universität Karlsruhe, 1996.